# Transformations for Direct Design of 2-D Filters from Appropriate 1-D Functions 

NIKOS E. MASTORAKIS<br>WSEAS Research Department<br>Agiou Ioannou Theologou 17-23<br>Zografou, 15773, Athens, GREECE<br>and<br>Technical University of Sofia, English Language Faculty<br>Industrial Engineering Department<br>Sofia 1000, BULGARIA<br>mastor@wseas.org<br>and with the Military Institutes of University Education (ASEI), Hellenic Naval Academy, Terma Hatzikyriakou, Piraeus, GREECE

Abstract: - In this paper, general transformations for designing 2-D (Two-Dimensional) IIR filters are provided. The present approach can be viewed as an extension of the classical method of magnitude approximation and can be applied in several cases of 2-D filter design. Numerical examples illustrate the validity and the efficiency of the method.

Key-Words: - 2-D Filters, FIR Filters, IIR Filters, Multidimensional Systems, Multidimensional Filters, Filter Design, McClellan Transformations

## 1 Introduction

Several techniques for the design of 2-D (TwoDimensional) filters have been published recently due to the vast applications of 2-D filters in image processing, satellite communications, biomedical imaging, computer vision etc.. Various popular methods have received considerable attention by engineers and research scholars. In general the design of 2-D FIR filters includes a Fourier method that uses Fourier analysis, where appropriate window Functions can also eliminate Gibbs' oscillations, a Transformations' method which is based on McClellan Transformations from appropriate 1-D filters [1],[2] and an optimization method i.e. the minimization of an appropriate norm, [1],[2]. Similarly, the design of 2-D IIR filters includes also transformations, Mirror Image Polynomials, SVD (Singular Value Decomposition) and Optimization, [1],[2]. Several Authors have published works on optimization-based 2-D filter design while a great number of papers are dedicated to transformations and mainly to McClellan Transformations, $\quad[3] \div[18] . \quad$ McClellan

Transformations were introduced in [3] and have been used for the last forty years in many theoretical topics and engineering applications. Harn and Shenoi pointed out in [5] and Nguyen and Swamy reported in [6], that till now a transformation for IIR filter design analogous to McClellan transformation does not exist due to the requirements of 2-D stability.
The purpose of this paper is to find such transformations and an attempt is made in section II.

This paper examines this transformation as well as its generalization to the general 2-D filters design. The usefulness of the proposed transforms is verified through two examples in section III. Finally, there is a conclusion.

## 2 The Transformation and its Generalizations

Let us consider the 1-D function of the prototype low-pass Butterworth filter (cut-off frequency $\Omega=1$ )

$$
\begin{equation*}
y=|H(j \Omega)|^{2}=\frac{1}{1+\Omega^{2 n}} \tag{1}
\end{equation*}
$$

The question here is "could it be possible after some appropriate transformation, for this function to derive the stable transfer function of a 2-D lowpass filter without the usual Transformations (McClellan Transformation) or the popular optimization techniques?"

An easy transformation could be for example $\Omega=\omega_{1}^{2}+\omega_{2}^{2}$, where $\omega_{1}, \omega_{2}$ are the frequencies of the discrete 2-D filter. However, in such a case, a filter with magnitude $\left|H\left(j \omega_{1}, j \omega_{2}\right)\right|=\frac{1}{1+\left(\omega_{1}^{2}+\omega_{2}^{2}\right)^{2 n}}$ can not be implemented since it does not correspond to a transfer function that would be a rational function of $z_{1}^{-1}, z_{2}^{-1}$. Another transformation could be, for instance, $\Omega=\sin ^{2}\left(\omega_{1}\right)+\sin ^{2}\left(\omega_{2}\right)$. The problem in this case is that
due to $\sin ^{2}(\varphi)=\frac{1-\cos (2 \varphi)}{2}$ the period of $\omega_{1}$ (and $\omega_{2}$ ) would be $\pi$ and not $2 \pi$ as the 2-D filter design demands.

The transformation $\Omega=\sin ^{2}\left(\frac{\omega_{1}}{2}\right)+\sin ^{2}\left(\frac{\omega_{2}}{2}\right)$ or the more general $\Omega=\sin ^{2 p}\left(\frac{\omega_{1}}{2}\right)+\sin ^{2 p}\left(\frac{\omega_{2}}{2}\right)$ seem more logical, where p is a positive integer.

Obviously, the low frequencies of $\Omega$ are depicted to a region of low frequencies of $\omega_{1}, \omega_{2}$-details will be presented in the next paragraphs.

In all the previous cases, the 2-D filter will be a noncausal 2-D system. The non-causality in 2-D systems is not a problem, since the dimensions are spatial and do not correspond to time.

By introducing $\Omega=\sin ^{2 p}\left(\frac{\omega_{1}}{2}\right)+\sin ^{2 p}\left(\frac{\omega_{2}}{2}\right)$ from our prototype function of (1) we take

$$
\begin{equation*}
\left|H\left(j \omega_{1}, j \omega_{2}\right)\right|=\frac{1}{1+\left(\sin ^{2 p}\left(\frac{\omega_{1}}{2}\right)+\sin ^{2 p}\left(\frac{\omega_{2}}{2}\right)\right)^{2 n}} \tag{2}
\end{equation*}
$$

So, using $\sin ^{2}(\varphi)=\frac{1-\cos (2 \varphi)}{2}$, one gets

$$
\left|H\left(j \omega_{1}, j \omega_{2}\right)\right|=\frac{1}{1+\left(\left(\frac{1-\cos \left(\omega_{1}\right)}{2}\right)^{p}+\left(\frac{1-\cos \left(\omega_{1}\right)}{2}\right)^{p}\right)^{2 n}}
$$

Considering now that our filter is of zero-phase, one can write
$H\left(j \omega_{1}, j \omega_{2}\right)=\frac{1}{1+\left(\left(\frac{1-\cos \left(\omega_{1}\right)}{2}\right)^{p}+\left(\frac{1-\cos \left(\omega_{1}\right)}{2}\right)^{p}\right)^{2 n}}$
that can be implemented because
$\cos \left(\omega_{1}\right)=\frac{z_{1}^{-1}+\left(z_{1}^{-1}\right)^{-1}}{2}, \cos \left(\omega_{2}\right)=\frac{z_{2}^{-1}+\left(z_{2}^{-1}\right)^{-1}}{2}$

Therefore
$H\left(z_{1}^{-1}, z_{2}^{-1}\right)=\frac{1}{1+\left(\left(\frac{2-z_{1}^{-1}-\left(z_{1}^{-1}\right)^{-1}}{4}\right)^{p}+\left(\frac{2-z_{2}^{-1}-\left(z_{2}^{-1}\right)^{-1}}{4}\right)^{p}\right)^{2 n}}$
which can easily implemented provided that this is a BIBO (Bounded Input Bounded Output Filter). BIBO Stability of (3) can be proven easily taking into account that we have a non-causal 2-D system. So, a necessary and sufficient condition is ([28)]
$B\left(z_{1}^{-1}, z_{2}^{-1}\right) \neq 0$ for all $z_{1}^{-1}, z_{2}^{-1}$ with $\left|z_{1}^{-1}\right|=1,\left|z_{2}^{-1}\right|=1$
where
$B\left(z_{1}^{-1}, z_{2}^{-1}\right)=1+\left(\left(\frac{2-z_{1}^{-1}-\left(z_{1}^{-1}\right)^{-1}}{4}\right)^{p}+\left(\frac{2-z_{2}^{-1}-\left(z_{2}^{-1}\right)^{-1}}{4}\right)^{p}\right)^{2 n}$

However, this can be proven easily, since $z_{1}^{-1}=e^{i \vartheta_{1}}$ and $z_{2}^{-1}=\rho e^{i \vartheta_{2}}$.

So, $B\left(z_{1}^{-1}, z_{2}^{-1}\right)=1+\left(\cos ^{2 p}\left(\vartheta_{1} / 2\right)+\cos ^{2 p}\left(\vartheta_{2} / 2\right)\right)^{2 n} \neq 0$
which is obviously $\neq 0$

Therefore the BIBO Stability of the filter in (3) has been proved.

A simple extension of (1) can be

$$
\begin{equation*}
y=|H(j \Omega)|^{2}=\frac{1}{1+\varepsilon^{2 n} \Omega^{2 n}} \tag{4}
\end{equation*}
$$

where $\varepsilon>0$ and cut-off frequency $\varepsilon^{-1}$.

Under the transformation $\Omega=\sin ^{2 p}\left(\frac{\omega_{1}}{2}\right)+\sin ^{2 p}\left(\frac{\omega_{2}}{2}\right)$, and following the same steps, one takes the stable 2D filter

$$
H\left(z_{1}^{-1}, z_{2}^{-1}\right)=\frac{1}{1+\varepsilon^{2 n}\left(\left(\frac{z_{1}^{-1}+\left(z_{1}^{-1}\right)^{-1}}{2}\right)^{p}+\left(\frac{z_{2}^{-1}+\left(z_{2}^{-1}\right)^{-1}}{2}\right)^{p}\right)^{2 n}}
$$

## 3 Numerical Examples and Design

## Example III. 1

Consider (1) with $n=8$, the magnitude $y=\frac{1}{1+\Omega^{16}}$ is depicted in Fig. 1


Fig. 1 Plot $y=\frac{1}{1+\Omega^{16}}$

Then, with the transformation $\Omega=\sin ^{4}\left(\frac{\omega_{1}}{2}\right)+\sin ^{4}\left(\frac{\omega_{2}}{2}\right)$, i.e. $p=2$, one gets the zerophase filter with magnitude response in Fig. 2


Fig.2. Plot $\left|H\left(j \omega_{1}, j \omega_{2}\right)\right|=\frac{1}{1+\left(\sin ^{4}\left(\frac{\omega_{1}}{2}\right)+\sin ^{4}\left(\frac{\omega_{2}}{2}\right)\right)^{16}}$

## Example III. 2

Consider again (1), with $n=8, y=\frac{1}{1+\Omega^{16}}$

Under the transformation $\Omega=\sin ^{8}\left(\frac{\omega_{1}}{2}\right)+\sin ^{8}\left(\frac{\omega_{2}}{2}\right)$, i.e. $p=4$, one gets the zero-phase filter with magnitude response in Fig. 3


Fig.3. Plot $H\left(j \omega_{1}, j \omega_{2}\right)=\frac{1}{1+\left(\sin ^{8}\left(\frac{\omega_{1}}{2}\right)+\sin ^{8}\left(\frac{\omega_{2}}{2}\right)\right)^{16}}$

By considering $\varepsilon=2$, we have $y=\frac{1}{1+\varepsilon^{16} \Omega^{16}}$

Now the transformation $\Omega=\sin ^{8}\left(\frac{\omega_{1}}{2}\right)+\sin ^{8}\left(\frac{\omega_{2}}{2}\right)$
$y=\frac{1}{1+\varepsilon^{16}\left(\sin ^{8}\left(\frac{\omega_{1}}{2}\right)+\sin ^{8}\left(\frac{\omega_{2}}{2}\right)\right)^{16}}$ yields
$\varepsilon^{16}\left(\sin ^{8}\left(\frac{\omega_{1}}{2}\right)+\sin ^{8}\left(\frac{\omega_{2}}{2}\right)\right)^{16}=1$
or equivalently


Fig. 4.
Plot $H\left(j \omega_{1}, j \omega_{2}\right)=\frac{1}{1+2^{16}\left(\sin ^{8}\left(\frac{\omega_{1}}{2}\right)+\sin ^{8}\left(\frac{\omega_{2}}{2}\right)\right)^{16}}$
with cut-off frequencies to be given by
$\sin ^{8}\left(\frac{\omega_{1}}{2}\right)+\sin ^{8}\left(\frac{\omega_{2}}{2}\right)=\frac{1}{2}$
or $\left\|\left(\sin \left(\frac{\omega_{1}}{2}\right), \sin \left(\frac{\omega_{2}}{2}\right)\right)\right\|_{8}=\sqrt[0.8]{.5} \approx 0.917$
On the other hand
$\left\|\left(\sin \left(\frac{\omega_{1}}{2}\right), \sin \left(\frac{\omega_{2}}{2}\right)\right)\right\|_{8} \approx \max \left(\left|\sin \left(\frac{\omega_{1}}{2}\right)\right|,\left|\sin \left(\frac{\omega_{2}}{2}\right)\right|\right)$
i.e. the cut-off frequencies can be defined by
$\omega_{1}, \omega_{2}=2 \cdot \arcsin (0.917)=2.3209$

Solution: Our filter must be

$$
y=\frac{1}{1+\varepsilon^{16}\left(a^{8} \sin ^{8}\left(\frac{\omega_{1}}{2}\right)+b^{8} \sin ^{8}\left(\frac{\omega_{2}}{2}\right)\right)^{16}}
$$

Without loss of generality the constant $\varepsilon$ is incorporated to $a$ and $b$. So, one can set $\varepsilon=1$. The cut-off frequencies are found from the equation
$\left(a^{8} \sin ^{8}\left(\frac{\omega_{1}}{2}\right)+b^{8} \sin ^{8}\left(\frac{\omega_{2}}{2}\right)\right)^{16}=1$
Or equivalently

$$
a^{8} \sin ^{8}\left(\frac{\omega_{1}}{2}\right)+b^{8} \sin ^{8}\left(\frac{\omega_{2}}{2}\right)=1
$$

Since,
$\left\|\left(a \sin \left(\frac{\omega_{1}}{2}\right), b \sin \left(\frac{\omega_{2}}{2}\right)\right)\right\|_{8} \approx \max \left(\left|a \sin \left(\frac{\omega_{1}}{2}\right)\right|,\left|b \sin \left(\frac{\omega_{2}}{2}\right)\right|\right)$
Using $\quad \omega_{10}= \pm \pi / 2, \quad \omega_{20}= \pm \pi / 3 \quad$ from $\max \left(\left|a \sin \left(\frac{\omega_{10}}{2}\right)\right|,\left|b \sin \left(\frac{\omega_{20}}{2}\right)\right|\right)=1$ we take
$a=\sqrt{2}, b=2$

Hence, our filter is
$H\left(z_{1}^{-1}, z_{2}^{-1}\right)=\frac{1}{1+\left(2^{4}\left(\frac{2-z_{1}^{-1}-\left(z_{1}^{-1}\right)^{-1}}{4}\right)^{4}+2^{8}\left(\frac{2-z_{2}^{-1}-\left(z_{2}^{-1}\right)^{-1}}{4}\right)^{4}\right)^{16}}$
and can be depicted easily in Fig. 5 taking
$\left|H\left(\omega_{1}, \omega_{2}\right)\right|=\frac{1}{1+\left(2^{4} \sin ^{8}\left(\frac{\omega_{1}}{2}\right)+2^{8} \sin ^{8}\left(\frac{\omega_{2}}{2}\right)\right)^{16}}$


Fig.5: Magnitude response of the filter of Example III. 3

In Fig.5, the magnitude response is illustrated if $\omega_{2}=0$


Fig.6: Magnitude response of the filter of Example III. 3 if $\omega_{2}=0$

## 4 Conclusion

In this brief, new transformations for designing 2D (Two-Dimensional) IIR filters are introduced. The resulting non-causal IIR filter is proven to be BIBO stable. Some numerical examples illustrate the validity and usefulness of the proposed transformation. Analogous transformations can be
derived for highpass, bandstop and bandpass noncausal 2-D IIR filters.

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