# A Method of Order $1+\sqrt{3}$ for Computing the Smallest Eigenvalue of a Symmetric Toeplitz Matrix 

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#### Abstract

In this note we discuss a method of order $1+\sqrt{3}$ for computing the smallest eigenvalue $\lambda_{1}$ of a symmetric and positive definite Toeplitz matrix. It generalizes and improves a method introduced in [7] which is based on rational Hermitean interpolation of the secular equation. Taking advantage of a further rational approximation of the secular equation which is essentially for free and which yields lower bounds of $\lambda_{1}$ we obtain an improved stopping criterion.


Keywords: eigenvalue problem, Toeplitz matrix, secular equation

## 1 Introduction

The problem of finding the smallest eigenvalue $\lambda_{1}$ of a real symmetric, positive definite Toeplitz matrix $T$ is of considerable interest in signal processing. Given the covariance sequence of the observed data, Pisarenko [11] suggested a method which determines the sinusoidal frequencies from the eigenvector of the covariance matrix associated with its minimum eigenvalue. The computation of the minimum eigenvalue of $T$ was studied in, e.g. [1], [4], [5], [6], [7], [8], [9], [10], [12], [13], [14].

In their seminal paper [1] Cybenko and Van Loan presented the following method: By bisection they first determine an initial approximation $\mu_{0} \in\left(\lambda_{1}, \omega_{1}\right)$ where $\omega_{1}$ denotes the smallest pole of the secular equation $f$, and they improve $\mu_{0}$ by Newton's method for $f$ which converges monotonely and quadratically to $\lambda_{1}$. By replacing Newton's method by a root finding method based on Rational Hermitean interpolation of $f$ Mackens and the second author in [7] improved this approach substantially.

In this note we revisit this method. In [7] the $k$-th iterate $\mu_{k}$ was chosen to be the unique root of

$$
g(\lambda)=a_{0}+a_{1}(\lambda-\alpha)+(\lambda-\alpha)^{2} \frac{b}{c-\lambda}
$$

in $\left(\alpha, \mu_{k-1}\right)$ where $\alpha$ is a lower bound of $\lambda_{1}$ obtained in the bisection phase, and $a_{0}, a_{1}, b$ and $c$ are chosen such that $g$ interpolates $f$ at $\alpha$ and $\mu_{k-1}$ in the Hermitean sense. It was proved that this method converges monotonely and quadratically to $\lambda_{1}$ and that it converges faster than Newton's method, i.e. if $\mu \in\left(\lambda_{1}, \omega_{1}\right)$ then the smallest root of $g$ is closer to $\lambda_{1}$ than the Newton iterate with initial guess $\mu$.

The method suffers the same disadvantage as the method of false position for convex or concave functions: one interpolation knot (in our case $\alpha$ ) is stationary, and only the other one converges momotonely to the wanted solution. In the root finding case one gains a substantial improvement if one drops the requirement that $f$ has opposite signs at the two interpolation knots and replaces the method of false position by the secant method. In this note we prove that the method in [7] can be improved in a similar way if one chooses the new iterate $\mu_{k}$ as the unique root of $g$ were the parameters $a_{0}, a_{1}, b$ and $c$ are determined such that $g$ and $g^{\prime}$ interpolate $f$ and $f^{\prime}$, respectively, at $\mu_{k}$ and $\mu_{k-1}$. It is shown that the order of convergence of this modified method is $1+\sqrt{3}$.

In [7] we based a stopping criterion on a lower bounds of $\lambda_{1}$ which are determined from a quadratic interpolation. This one is improved using a further rational interpolation of $f$ with a fixed pole which is obtained for free in the course of the algorithm.

## 2 Rational Hermitean interpolation

Let $T \in \mathbb{R}^{(n, n)}$ be a symmetric positive definite Toeplitz matrix. We assume that its diagonal is normalized and consider the following partition:

$$
T=\left(\begin{array}{cc}
1 & t^{T} \\
t & G
\end{array}\right)
$$

It is well known that the eigenvalues of $T$ and of $G$ are real and positive and satisfy the interlacing property $\lambda_{1} \leq \omega_{1} \leq \lambda_{2} \leq \ldots \leq \omega_{n-1} \leq \lambda_{n}$ where $\lambda_{j}$ and $\omega_{j}$ is the $j$ th smallest eigenvalue of $T$ and its principal submatrix $G$, respectively.

We assume that $\lambda_{1}<\omega_{1}$. Then $\lambda_{1}$ is the smallest root of the secular equation

$$
\begin{equation*}
f(\lambda):=-1+\lambda+t^{T}(G-\lambda I)^{-1} t=0 \tag{1}
\end{equation*}
$$

It is easily seen that $f$ is strictly monotonely increasing and strictly convex in the interval $\left(0, \omega_{1}\right)$, and therefore for every initial value $\mu_{0} \in\left(\lambda_{1}, \omega_{1}\right)$ Newton's method converges monotonely decreasing and quadratically to $\lambda_{1}$.

Cybenko and Van Loan [1] suggested to determine an initial value $\mu_{0}$ by bisection based on Durbin's algorithm (cf. [2], p. 184 ff ). If $\mu$ is not in the spectrum of any of the principal submatrices of $T-\mu I$ then Durbin's algorithm applied to $(T-\mu I) /(1-\mu)$ determines a lower triangular matrix

$$
L=\left(\begin{array}{cccc}
1 & 0 & \ldots & 0 \\
\ell_{21} & 1 & \ldots & 0 \\
\vdots & \vdots & \ddots & \\
\ell_{n 1} & \ell_{n 2} & \ldots & 1
\end{array}\right)
$$

such that

$$
\begin{equation*}
\frac{1}{1-\mu} L(T-\mu I) L^{T}=D:=\operatorname{diag}\left\{1, E_{1}, \ldots, E_{n-1}\right\} \tag{2}
\end{equation*}
$$

If $\tilde{L}$ is obtained from $L$ by dropping the last row and last column then obviously

$$
\frac{1}{1-\mu} \tilde{L}(G-\mu I) \tilde{L}^{T}=\tilde{D}:=\operatorname{diag}\left\{1, E_{1}, \ldots, E_{n-2}\right\}
$$

Hence, from Sylvester's law of inertia one gets
(i) $\mu<\lambda_{1}$, if $E_{j}>0$ for $j=1, \ldots, n-1$,
(ii) $\mu \in\left[\lambda_{1}, \omega_{1}\right)$, if $E_{j}>0$ for $j=1, \ldots, n-2$ and $E_{n-1} \leq 0$,
(iii) and $\mu>\omega_{1}$, if $E_{j}<0$ for some $j \in\{1, \ldots, n-$ $2\}$.

An upper bound of $\lambda_{1}$ to start the bisection process can be obtained in the following way. Let $w:=-G^{-1} t$ be the solution of the Yule-Walker system. Then

$$
q:=\frac{1}{1+t^{T} w}\binom{1}{w}=T^{-1} e^{1}
$$

is the first iterate of the inverse iteration with shift parameter 0 starting with the unit vector $e_{1}$ which can be expected to be not too bad an approximation of the eigenvector corresponding to the smallest eigenvalue $\lambda_{1}$. The Rayleigh quotient

$$
\begin{equation*}
R(q):=\frac{q^{T} T q}{q^{T} q}=\frac{1+t^{T} w}{1+\|w\|_{2}^{2}} \tag{3}
\end{equation*}
$$

is an upper bound of $\lambda_{1}$ which should be not too bad either.

Since

$$
\begin{equation*}
f^{\prime}(\lambda)=1+\left\|(G-\lambda I)^{-1} t\right\|_{2}^{2} \tag{4}
\end{equation*}
$$

a Newton step can be performed in the following way:

$$
\begin{aligned}
& \text { Solve }\left(G-\mu_{k} I\right) w=-t \quad \text { for } \quad w \\
& \text { and set } \mu_{k+1}:=\mu_{k}-\frac{-1+\mu_{k}-w^{T} t}{1+\|w\|_{2}^{2}}
\end{aligned}
$$

where the Yule-Walker system

$$
\begin{equation*}
(G-\mu I) w=-t \tag{5}
\end{equation*}
$$

can be solved by Durbin's algorithm requiring $2 n^{2}$ flops.

The global convergence behaviour of Newton's method usually is not satisfactory since the smallest root $\lambda_{1}$ and the smallest pole $\omega_{1}$ of the rational function $f$ can be very close to each other. In this situation the initial steps of Newton's method are extremely slow, at least if the initial guess is close to $\omega_{1}$.

Approximating the secular equation by a suitable rational function the convergence of the method (i.e. the bisection phase and the root finding by

Newton's method) can be improved considerably. In terms of condensation methods (cf. [3]) the secular equation $f$ can be interpreted as the exact condensation of the eigenvalue problem $T x=\lambda x$ where $x_{2}, \ldots, x_{n}$ are chosen to be slaves and $x_{1}$ is the only master. Using spectral information of the slave problem $(G-\mu I) v=0$ the function $f$ obtains the form (cf. [3])

$$
f(\lambda)=f(0)+f^{\prime}(0) \lambda+\lambda^{2} \sum_{j=1}^{n-1} \frac{\alpha_{j}^{2}}{\omega_{j}-\lambda}
$$

where $\alpha_{j}, j=1, \ldots, n-1$, are real numbers depending on the eigenvectors of $G$. With a shift $\mu$ which is not in the spectrum of $G f$ can be rewritten as

$$
\begin{equation*}
f(\lambda)=f(\mu)+(\lambda-\mu) f^{\prime}(\mu)+(\lambda-\mu)^{2} \phi(\lambda ; \mu) \tag{6}
\end{equation*}
$$

where

$$
\begin{equation*}
\phi(\lambda ; \mu)=\sum_{j=1}^{n-1} \frac{\alpha_{j}^{2} \gamma_{j}^{2}}{\omega_{j}-\lambda}, \quad \gamma_{j}=\frac{\omega_{j}}{\omega_{j}-\mu} \tag{7}
\end{equation*}
$$

The representation (6) and (7) of $f$ suggests to replace the linearization of $f$ in Newton's method by a root finding method based on a rational model

$$
\begin{equation*}
g(\lambda ; \mu, \nu)=f(\mu)+(\lambda-\mu) f^{\prime}(\mu)+(\lambda-\mu)^{2} \frac{b}{c-\lambda} \tag{8}
\end{equation*}
$$

where $\mu$ and $\nu$ are given approximations to $\lambda_{1}$ and $b$ and $c$ are determined such that

$$
\begin{equation*}
g(\nu ; \mu, \nu)=f(\nu), g^{\prime}(\nu ; \mu, \nu)=f^{\prime}(\nu) \tag{9}
\end{equation*}
$$

Theorem 1: Let $g$ be given by (8) and (9) where $\mu$ and $\nu$ are not in the spectrum of $G$. Then

$$
\begin{equation*}
b=\frac{\phi(\nu ; \mu)^{2}}{\phi^{\prime}(\nu ; \mu)} \geq 0, c=\nu+\frac{\phi(\nu ; \mu)}{\phi^{\prime}(\nu ; \mu)} \geq \omega_{1} \tag{10}
\end{equation*}
$$

Proof: From equations (6) and (8) we obtain

$$
\begin{equation*}
g(\lambda ; \mu, \nu)-f(\lambda)=(\lambda-\mu)^{2}\left(\frac{b}{c-\lambda}-\phi(\lambda ; \mu)\right) . \tag{11}
\end{equation*}
$$

Hence the interpolation conditions (9) yield

$$
\frac{b}{c-\nu}-\phi(\nu ; \mu)=0, \frac{b}{(c-\nu)^{2}}-\phi^{\prime}(\nu ; \mu)=0
$$

from which we get the representations of $b$ and $c$ in (10).
$b \geq 0$ is obvious, and $c \geq \omega_{1}$ follows from

$$
c=\sum_{j=1}^{n-1} \frac{\alpha_{j}^{2} \gamma_{j}^{2}}{\left(\omega_{j}-\nu\right)^{2}} \omega_{j} / \sum_{j=1}^{n-1} \frac{\alpha_{j}^{2} \gamma_{j}^{2}}{\left(\omega_{j}-\nu\right)^{2}}
$$

which is obtained from (7) and (10).
Theorem 2: If $\mu$ and $\nu$ are not in the spectrum of $G$ it holds

$$
\begin{equation*}
f(\lambda)-g(\lambda)=(\lambda-\mu)^{2}(\lambda-\nu)^{2} \psi(\lambda ; \mu, \nu) \tag{12}
\end{equation*}
$$

where $\psi=\psi_{1} / \psi_{2}$,

$$
\begin{gathered}
\psi_{1}=\sum_{1 \leq j<k \leq n-1} \frac{\alpha_{j}^{2} \alpha_{k}^{2} \omega_{j}^{2} \omega_{k}^{2}\left(\omega_{k}-\omega_{j}\right)^{2}}{\tau_{j k}(\mu)^{2} \tau_{j k}(\nu)^{2}\left(\omega_{j}-\lambda\right)\left(\omega_{k}-\lambda\right)} \\
\tau_{j k}(\lambda)=\left(\omega_{j}-\lambda\right)\left(\omega_{k}-\lambda\right)
\end{gathered}
$$

and

$$
\psi_{2}=\sum_{j=1}^{n-1} \frac{\alpha_{j}^{2} \omega_{j}^{2}}{\left(\omega_{j}-\mu\right)^{2}\left(\omega_{j}-\nu\right)^{2}}\left(\omega_{j}-\lambda\right)
$$

Proof: From equations (10) and (11) it follows

$$
\begin{aligned}
& f(\lambda)-g(\lambda)= \\
& \quad(\lambda-\mu)^{2}\left(\phi(\lambda ; \mu)-\frac{\phi(\nu ; \mu)^{2}}{\phi(\nu ; \mu)+(\nu-\lambda) \phi^{\prime}(\nu ; \mu)}\right),
\end{aligned}
$$

and taking advantage of (7) an easy but lengthy calculation yields (12).

In particular we obtain from Theorem $2 g\left(\lambda_{1}\right)<$ 0 , and since $g$ is strictly monotonely increasing and strictly convex in $[0, c) \supset\left[0, \omega_{1}\right)$ and $\lim _{\lambda \uparrow c} g(\lambda ; \mu, \nu)=\infty$ the unique root of $g$ in $[0, c)$ is an upper bound of the smallest eigenvalue $\lambda_{1}$ of $T$.

Assume that we are given a lower bound $\mu_{0}$ of $\lambda_{1}$ and an upper bound $\mu_{1} \in\left(\lambda_{1}, \omega_{1}\right)$ which is obtained by bisection, e.g. Then the unique root $\mu_{2}$ of $g\left(\cdot ; \mu_{1}, \mu_{0}\right)$ in $(0, c)$ satisfies $\lambda_{1} \leq \mu_{2}<\mu_{1}$. Mackens and the second author in [7] considered a method of false position like iteration where $\mu_{k+1}$ is defined as the unique root of $g\left(\cdot ; \mu_{k}, \mu_{0}\right)$, and they proved this method to be quadratically convergent.

Here we study the method which corresponds to the secant method where $\mu_{k+1}$ is determined as the unique root of $g\left(\cdot ; \mu_{k}, \mu_{k-1}\right)$. Again this algorithm yields a monotonely decreasing sequence $\left\{\mu_{k}\right\}$ which is bounded below by $\lambda_{1}$. The following Theorem 3 proves the convergence of this sequence to $\lambda_{1}$ and its order of convergence $1+\sqrt{3}$.
Theorem 3: Let $\mu_{1} \in\left(\lambda_{1}, \omega_{1}\right)$ and for $k \geq 2$ let $\mu_{k+1}$ be the unique root of $g\left(\cdot ; \mu_{k}, \mu_{k-1}\right)$ in $\left[0, \omega_{1}\right)$.

Then the sequence $\left\{\mu_{k}\right\}$ converges monotonely decreasing to $\lambda_{1}$, and its $R$-order of convergence is $1+\sqrt{3}$.
Proof: Let $\epsilon_{k}:=\mu_{k}-\lambda_{1}$. From $g\left(\mu_{k+1} ; \mu_{k}, \mu_{k-1}\right)=$ 0 and Theorem 2 we obtain for some $\xi_{k} \in\left(\lambda_{1}, \mu_{k+1}\right)$

$$
\begin{aligned}
& f\left(\mu_{k+1}\right)-f\left(\lambda_{1}\right)=f^{\prime}\left(\xi_{k}\right) \epsilon_{k+1}= \\
& \quad\left(\mu_{k}-\mu_{k+1}\right)^{2}\left(\mu_{k-1}-\mu_{k+1}\right)^{2} \psi\left(\mu_{k+1}, \mu_{k}, \mu_{k-1}\right)
\end{aligned}
$$

The sequence $\left\{\mu_{k}\right\}$ is monotonely decreasing and bounded away from $\omega_{1}$. Hence there exists $C>0$ such that

$$
\epsilon_{k+1} \leq C \epsilon_{k}^{2} \epsilon_{k-1}^{2}
$$

and for $e_{k}:=C^{1 / 3} \epsilon_{k}$ it holds

$$
e_{k+1} \leq e_{k}^{2} e_{k-1}^{2}
$$

Let $p=1+\sqrt{3}$ and $\eta:=\min \left(e_{0}, e_{1}^{1 / p}\right)$. We prove by induction

$$
\begin{equation*}
e_{k} \leq \eta^{\left(p^{k}\right)} \tag{13}
\end{equation*}
$$

which demonstrates that the R -order of convergence of $\mu_{k}$ equals $1+\sqrt{3}$.

For $k=0$ and $k=1$ (13) is trivial. If it hold for integers up to $k$ then it follows from $2(1+p)=p^{2}$

$$
\begin{aligned}
e_{k+1} & \leq e_{k}^{2} e_{k-1}^{2} \leq \eta^{\left(2 p^{k}\right)} \eta^{\left(2 p^{k-1}\right)} \\
& =\eta^{\left(2(1+p) p^{k-1}\right)}=\eta^{\left(p^{k+1}\right)}
\end{aligned}
$$

With a further rational interpolation of the secular equation we are able to construct a lower bound of $\lambda_{1}$. This will be the basis of our stopping criterion.
Theorem 4: Let $\kappa \in\left(0, \lambda_{1}\right), \mu \in\left(\kappa, \omega_{1}\right)$ and $p \in\left(\kappa, \omega_{1}\right)$. Let

$$
h(\lambda):=f(\mu)+f^{\prime}(\mu)(\lambda-\mu)+(\lambda-\mu)^{2} \frac{b}{p-\lambda}
$$

where $b$ is determined such that the interpolation condition $h(\kappa)=f(\kappa)$ holds.

Then $b>0$, i.e. $h$ is strictly monotonely increasing and strictly convex in $(0, p)$, and the unique root of $h$ in $(0, p)$ is a lower bound of $\lambda_{1}$.
Proof: From equation (6) and from the interpolation condition $h(\kappa)=f(\kappa)$ we obtain

$$
b=(p-\kappa) \phi(\kappa ; \mu)>0
$$

That the unique root $\tilde{\lambda}$ of $h$ in $(0, p)$ is a lower bound of $\lambda_{1}$ is obvious for $p \leq \lambda_{1}$. For $p>\lambda_{1}$ we have to
show $h\left(\lambda_{1}\right)>0$. This follows from equations (6) and (7):

$$
\begin{aligned}
& h\left(\lambda_{1}\right)=f(\mu)+f^{\prime}(\mu)\left(\lambda_{1}-\mu\right)+\left(\lambda_{1}-\mu\right)^{2} \frac{b}{p-\lambda_{1}} \\
& \quad=f\left(\lambda_{1}\right)-\left(\lambda_{1}-\mu\right)^{2}\left(\phi\left(\lambda_{1}\right)-\frac{(p-\kappa) \phi(\kappa)}{p-\lambda_{1}}\right) \\
& \quad=\frac{\left(\lambda_{1}-\mu\right)^{2}}{p-\lambda_{1}}\left((p-\kappa) \phi(\kappa)-\left(p-\lambda_{1}\right) \phi\left(\lambda_{1}\right)\right) \\
& \quad=\frac{\left(\lambda_{1}-\mu\right)^{2}}{p-\lambda_{1}} \sum_{j=1}^{n-1} \gamma_{j}^{2}\left(\frac{p-\kappa}{\omega_{j}-\kappa}-\frac{p-\lambda_{1}}{\omega_{j}-\lambda_{1}}\right) \\
& \quad=\frac{\left(\lambda_{1}-\mu\right)^{2}}{p-\lambda_{1}} \sum_{j=1}^{n-1} \gamma_{j}^{2} \frac{\left(\omega_{j}-p\right)\left(\lambda_{1}-\kappa\right)}{\left(\omega_{j}-\kappa\right)\left(\omega_{j}-\lambda_{1}\right)}>0
\end{aligned}
$$

Theorem 4 can be used to construct lower bounds of $\lambda_{1}$ in the course of the algorithm which are essentially for free. We already pointed out that Durbin's algorithm determines the factorization of $T-\mu I$ given in (2). Hence, solving the Yule-Walker system for some $\mu$ we can evaluate the characteristic polynomial

$$
\chi(\mu)=(1-\mu) E_{1} \cdot \ldots \cdot E_{n-2}
$$

of $G$ at negligible cost. Moreover, $\chi(\lambda)($ or $-\chi(\lambda))$ is monotonely decreasing and convex for $\lambda \leq \omega_{1}$. Therefore, if $\chi\left(\mu_{1}\right)$ and $\chi\left(\mu_{2}\right)$ are known for $\mu_{1}, \mu_{2} \in\left[0, \omega_{1}\right)$ then a secant step for $\chi$ yields an improved lower bound of $\omega_{1}$.

## 3 A MATLAB progam

The following MATLAB program determines a lower bound $\mu$ of the smallest eigenvalue of a symmetric and positive definite Toeplitz matrix which is given by the vector $t$ of dimension $n$. It uses the function [ $\mathrm{f}, \mathrm{Df}, \mathrm{chi}, \mathrm{loc}]=$ durbin $(\mathrm{mu}, \mathrm{t}, \mathrm{n})$ which returns the value f of the secular equation at $\mu$, its derivative Df , the value chi of the characteristic polynomial of $G$, and the location

$$
\text { loc }=\left\{\begin{array}{lll}
0 & \text { if } & \mu<\lambda_{1} \\
1 & \text { if } & \lambda_{1} \leq \mu<\omega_{1} \\
2 & \text { if } & \mu>\omega_{1}
\end{array}\right.
$$

of mu within the spectrum of $T$.
The functions rat_app and rat_app_fp return the smallest positive root of the rational function $g$ and $h$, respectively.

```
function mu=toeplitz_ev(t,n,tol)
[f0,Df0,chi0,loc]=durbin(0,t,n);
mu0=0; p=0;
lb=0; ub=-f0/Df0;
ka=0; fka=f0;
mu=rand*ub;
rel_err=1;
while abs(rel_err) > tol
    [f,Df,chi,loc]=durbin(mu,t,n);
    if loc == 2
        lambda=rat_app(mu0,f0,Df0,mu,f,Df);
        ub=min([mu,ub,lambda]);
        mu=0.5*(lb+ub);
    else
        p=max(p,mu-(mu-mu0)*chi/(chi-chi0));
        lb=rat_app_fp(ka,fka,mu,f,Df,p);
        root=rat_app(mu0,f0,Df0,mu,f,Df));
        ub=min(ub,root);
        mu0=mu;f0=f;Df0=Df;chi0=chi;
        if loc == 0
            ka=mu0;fka=f0;end
        rel_err=ub/lb-1;
        if loc == 1
            mu=root;
        else
            newt=mu-f0/Df0;
            if abs((root-newt)/root)<0.01
                mu=root;
            else
                mu=0.1*lb+0.9*ub;
                    end
            end
        end
    end
```

Some remarks are in order.
$2-3: \mathrm{mu} 0<\lambda_{1}$ with known $\mathrm{f} 0=f(\mathrm{mu} 0)$, $\operatorname{Df} 0=f^{\prime}(\mathrm{mu} 0)$ and chi0 $=\chi(\mathrm{muO})$ is one knot in the rational interpolation of $f$ and the secant method for $\chi$. p is a lower bound of $\omega_{1}$ used to determine a lower bound of $\lambda_{1}$.

4 : lb is a lower bound of $\lambda_{1}$ and ub an upper bound. $u b=-s 0 / D s 0$ is obtained from (6).

5 : ka is a lower bound of $\lambda_{1}$ with known $\mathrm{fka}=f(\mathrm{ka})$ which corresponds to $\kappa$ in Theorem 4.

6 : The algorithm starts with a test parameter mu randomly chosen in the interval [lb, ub].
$10-13$ : By Theorem 2 the smallest root lambda of $g(\cdot ; \mathrm{mu}, \mathrm{muO})$ is an upper bound of $\lambda_{1}$. It is for free, and in some cases it is smaller than mu. This modification of the bisection method actually improves the performance of the method.

15 : The lower bound p of the pole might be improved by a secant step for the characteristic polynomial of $G$.
$16: \mathrm{lb}$ is the lower bound of $\lambda_{1}$ from Theorem 4.

17-18: The root of $g(\cdot ; \mathrm{mu}, \mathrm{mu} 0)$ is an upper bound of $\lambda_{1}$, and it further enhances the bisection method.
$20-21$ : If mu< $\lambda_{1}$, mu can be used as $\kappa$ of Theorem 4 in subsequent iteration steps.
$23-24$ : For $m u \in\left(\lambda_{1}, \omega_{1}\right)$ the method continues with test parameter mu=root.
$25-32:$ For $\mathrm{mu}<\lambda_{1}$ we introduce a tie break rule which was motivated in [7]. newt is the result of a Newton step for $f$. Hence root and newt are second order approximations of $\lambda_{1}$. If they are not close to each other the test parameter mu can not be close to $\lambda_{1}$. In this case we reduce the next test parameter. This modification improves the performance of the method, in particular if the gap between $\lambda_{1}$ and $\omega_{1}$ is very narrow.

## 4 Numerical results

To test the method we considered the following class of Toeplitz matrices:

$$
\begin{equation*}
T=m \sum_{k=1}^{n} \eta_{k} T_{2 \pi \theta_{k}} \tag{14}
\end{equation*}
$$

where $m$ is chosen such that the diagonal of $T$ is normalized to $t_{0}=1$,

$$
T_{\theta}=\left(T_{i j}\right)=(\cos (\theta(i-j)))
$$

and $\eta_{k}$ and $\theta_{k}$ are uniformly distributed random numbers in the interval $[0,1]$ (cf. Cybenko and Van Loan [1]).

Table 1 contains the average number of flops and the average number of Durbin steps needed to determine the smallest eigenvalue in 100 test problems with each of the dimensions $n=32,64,128,256$, 512,1024 and 2048. The iteration was terminated if the error was guaranteed to be less than $10^{-6}$ by the error bound from Theorem 4. For comparison we added the results for the quadratically convergent method in [7].

| dim. | order $1+\sqrt{3}$ |  | method in [7] |  |
| :---: | :---: | :---: | :---: | :---: |
|  | flops | steps | flops | steps |
| 32 | 1.086 E04 | 4.34 | 1.153 E04 | 4.67 |
| 64 | 4.639 E04 | 5.14 | 4.669 E04 | 5.39 |
| 128 | 1.804 E05 | 5.25 | 1.900 E05 | 5.79 |
| 256 | 7.837 E05 | 5.84 | 8.790 E05 | 6.85 |
| 512 | 3.512 E06 | 6.62 | 3.892 E06 | 7.69 |
| 1024 | 1.531 E 07 | 7.26 | 1.730 E07 | 8.75 |
| 2048 | 6.268 E07 | 7.45 | 7.590 E07 | 9.59 |

Tab. 1.

## 5 Concluding Remarks

We have presented an algorithm for computing the smallest eigenvalue of a symmetric and positive definite Toeplitz matrix of order $1+\sqrt{3}$. Realistic error bounds were obtained at negligible cost. We used Durbin's algorithm to solve the occuring YuleWalker systems and to determine the Schur parameters $E_{j}$ requiring $2 n^{2}$ flops. This information can be gained from superfast Toeplitz solvers the complexity of which is only $O\left(n \log ^{2} n\right)$ operations. In a similar way as in [12] or [13] the method can be enhanced taking advantage of symmetry properties of the eigenvectors of $T$.

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