

# A chaotic oscillator using the Van der Pol dynamic immersed into a Jerk system

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## Abstract

In the present work a new chaotic oscillator is presented by immersing the dynamics of the Van der Pol oscillator into a Jerk system (a third order differential equation). The map of Poincaré was used to prove chaotic behavior.

## Key-words

Chaotic oscillator, Van der Pol oscillator, Jerk systems.

## I. INTRODUCTION

In the paper presented by J. C. Sprott in 2000 [1], Sprott shows many Jerk systems that exhibit chaotic behavior. A great part of these chaotic systems can be implemented using analog electronics ([1]). Also, a huge amount of chaotic oscillators have appeared recently (see, for instance, the references given in [1]). In this paper, we present a new chaotic oscillator obtained by immersing the dynamics of a slightly modified Van der Pol oscillator into a Jerk system. The chaotic test was done using the map of Poincaré (Banerjee, page 61, [3]).

## II. JERK SYSTEMS AND CHAOTIC OSCILLATORS

A Jerk system is a third order differential equation (a nonlinear one) of the form [1]:

$$\ddot{x} = f(\ddot{x}, \dot{x}, x) \quad (1)$$

where  $x, \dot{x}, \ddot{x}$ , and  $\ddot{x}$  represent position, velocity, acceleration and Jerk (the time derivative of the acceleration), respectively. For instance, the following Jerk system:

$$\ddot{x} + \ddot{x} - x\dot{x} + ax + b = 0 \quad (2)$$

presents chaotic motion with  $a=0.9$  and  $b=0.4$  [1]. Its strange attractor is shown in Fig. 1.

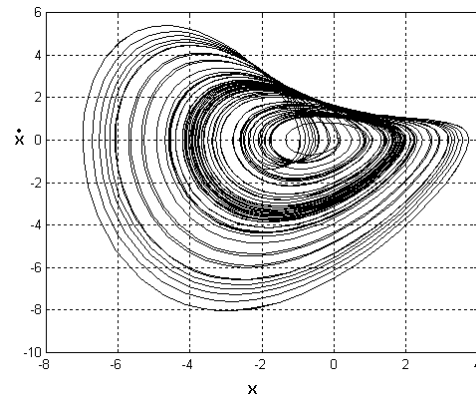


Fig. 1 Phase portrait using  $x(0) = \dot{x}(0) = \ddot{x}(0) = 0.1$ .

The Van der Pol dynamics is given by [2]:

$$\ddot{x} + \varepsilon(x^2 - 1)\dot{x} + x = 0 \quad (3)$$

which is a second order differential equation. Its phase portrait is pictured in Fig. 2 where it is shown the presence of a limit cycle (a stable one).

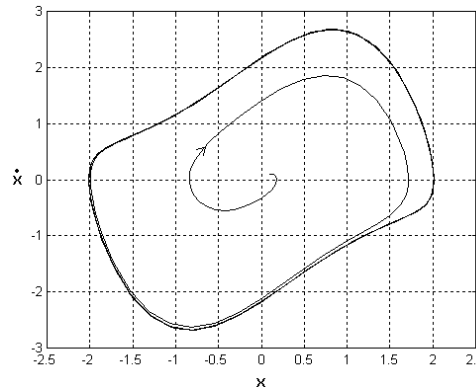


Fig. 2 Phase portrait using  $x(0) = \dot{x}(0) = 0.1$  and  $\varepsilon=1$ .

Motivated by the fact that a Jerk system can produce chaos, first, we slightly modify the dynamics of Van der Pol as follows:

$$a\ddot{x}^2 \text{sign}(\ddot{x}) + \varepsilon(x^2 - 1)\dot{x} + x = 0 \quad (4)$$

where  $a$  is a constant; after that, we immerse (4) into a Jerk system:

$$\begin{aligned} \ddot{x} &= f(\ddot{x}, \dot{x}, x) \\ &= -[a\ddot{x}^2 \text{sign}(\ddot{x}) + \varepsilon(x^2 - 1)\dot{x} + x] \end{aligned} \quad (5)$$

Its phase portrait (of (5)), with  $a=0.01$  and  $\varepsilon=7$ , is shown in Figure 3 and Figure 4, where  $x(0) = \dot{x}(0) = \ddot{x}(0) = 0.1$ .

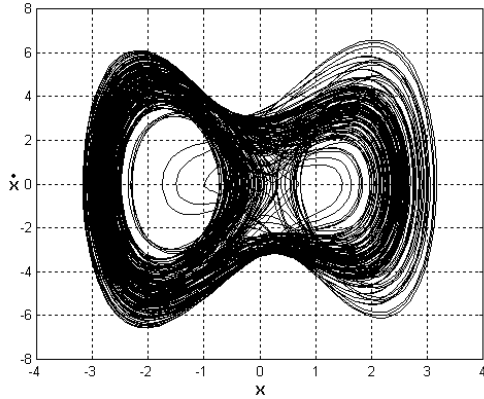


Fig. 3 Phase portrait.

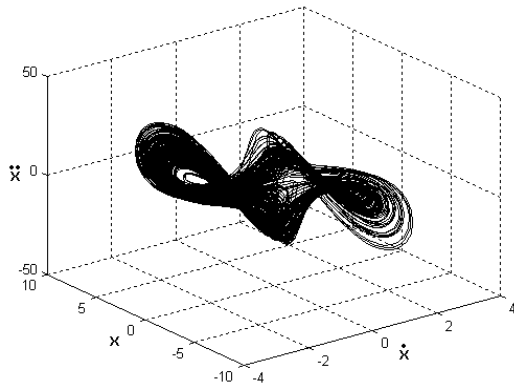


Fig. 4 Three-dimensional phase portrait.

Other values of  $a$  and  $\varepsilon$  that produce chaos are  $(a, \varepsilon) = \{(0.03, 7), (0.01, 6)\}$ . Hereafter we will use  $a=0.01$  and  $\varepsilon=7$ .

### III. CHAOTIC TEST

The chaotic test we use to prove chaotic motion for system (5) is the map of Poincaré

[3]. This map is obtained when we place, in appropriate form, a surface, called the "Poincaré section", into the state space [3] (see Fig. 5).

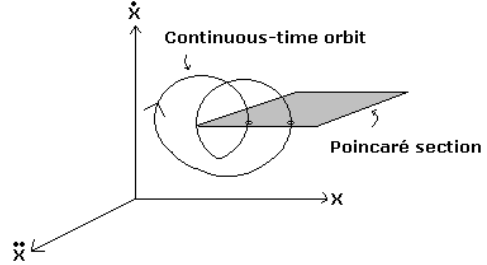


Fig. 5 Poincaré map.

In this way, and roughly speaking, the Poincaré map is a mapping of the points in which the trajectory intersects the Poincaré section. If the system is periodic, the Poincaré map will contain a few points; however, if the trajectory is chaotic, then, the Poincaré map will contain a huge amount of points distributed over the Poincaré section producing something similar to a strange attractor [3]. In Fig. 6 the map of Poincaré is shown placing the Poincaré section in the plane  $(x - \dot{x})$  in  $\ddot{x} = 0$  of our system.

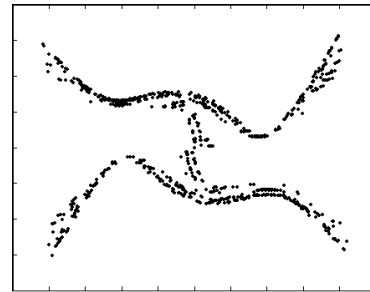


Fig. 6 Poincaré map.

From Fig. 6 follows that the proposed systems is a chaotic Jerk system.

### References

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