# A New Approach of the Induction Machine Mathematical Model 

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#### Abstract

In this paper an original approach based on the asymptotic approximation and the system theory was used to solve the nonlinear mathematical model of the induction machine. The mathematical model of induction machine has several nonlinear equations that cannot be integrated to determine the input/output equations or transfer function. The main objective of this paper is to find a method to solve the equations of mathematical model. For this reason first contribution was to replace speed with an asymptotic function then decupling the feedback loop and replacing it with an equivalent feedback control voltage. The second contribution consist in a special arrangement of mathematical model matrix to allow splitting off the decomposition model, consisting of forth order transition matrix, in two diagonal second order matrices that permit to write a linear state model. This interesting model was simulated and has the same state and output variation very close to the consecrated models. This approach permits to determine the solutions of model and the transfer matrices on the channels input (voltages), state (fluxes) and output (currents). At the final of the paper the theoretically results are validated by simulation using a real induction machine with good results.


Key-Words: - induction machine, special approximated model, state equations, solutions of equations transfer matrix, modeling, simulation

## 1 Introduction

Consider the normal three phase induction machine in the normal conditions of electro mechanic and electromagnetic parameters. We work in a vector space of continuous and derivable functions having two orthogonal Clarke coordinates d , q with s , r indexes for stator and respectively rotor.

The notations $\Psi_{\mathrm{i},}, \mathrm{i}_{\mathrm{ij}}, \mathrm{u}_{\mathrm{i}}$, represents the fluxes, currents and voltages, $L_{i}, \quad R_{i j}$, represent the inductances and resistances. The M, Ms, $\omega$, J are the torque, load, speed and inertia of drive [3].

The well known 10 machine equations from which three are nonlinear are given below:

$$
\begin{aligned}
& \dot{\Psi}_{s d}=-R_{s} \cdot i_{s d}+u_{s d} \\
& \dot{\Psi}_{s q}=-R_{s} \cdot i_{s q}+u_{s q} \\
& \dot{\Psi}_{r d}=-R_{r} \cdot i_{r d}+u_{r d}-\omega \cdot \Psi_{r q} \\
& \dot{\Psi}_{r q}=-R_{r} \cdot i_{s d}+u_{s d}+\omega \cdot \Psi_{r d} \\
& \Psi_{s d}=L_{s} \cdot i_{s d}+L_{m} \cdot i_{r d} \\
& \Psi_{s q}=L_{s} \cdot i_{s q}+L_{m} \cdot i_{r q} \\
& \Psi_{r d}=L_{r} \cdot i_{r d}+L_{m} \cdot i_{s d} \\
& \Psi_{r q}=L_{r} \cdot i_{r q}+L_{m} \cdot i_{s q}
\end{aligned}
$$

$$
\begin{aligned}
& \dot{\omega}=\frac{M-M_{s}}{J} \\
& M=\frac{3}{2} \cdot \frac{L_{m}}{L_{s} \cdot L_{r}-L_{m}^{2}}\left(\Psi_{r d} \cdot \Psi_{r q}-\Psi_{r q} \cdot \Psi_{s d}\right)
\end{aligned}
$$

This model cannot be integrated due to lack of nonlinear equations for the input/output transfer function. That is a difficult problem especially in control applications which require the identification of parameters, the estimation for state or the implementation of advanced algorithms [7]. Of course there are many solution which use approximations by linearization in steady state regime, but those are used in limited applications.

Now first we organize the two order vector space. Supply voltage vector on d-coordinate:
$u_{d a}=\left[u_{s d}, u_{r d}\right]^{T}=\left[U_{s m} \sin \omega_{0} \cdot t,-U_{r m} \cdot \cos \omega_{0} \cdot t\right]^{T}$
Supply voltage vector on q-coordinate:
$u_{q a}=\left[u_{s q}, u_{r q}\right]^{T}=\left[-U_{s m} \cos \omega_{0} \cdot t,-U_{r m} \cdot \sin \omega_{0} \cdot t\right]^{T}$
Fluxes are two components state vector:
$\Psi_{d}=\left[\Psi_{s d}, \Psi_{r d}\right]^{T}, \Psi_{q}=\left[\Psi_{s q}, \Psi_{r q}\right]^{T}$
Currents are two components output vectors:
$i_{d}=\left[i_{s d}, i_{r d}\right]^{T}, i_{q}=\left[i_{s q}, i_{r q}\right]^{T}$
Feedback voltage vector on d coordinate:
$u_{f q}=B \cdot \Psi_{q} \cdot \omega=\left[0, \omega \cdot \Psi_{r q}\right]^{T}$

Feedback voltage vector on q coordinate:
$u_{f d}=B \cdot \Psi_{d} \cdot \omega=\left[0, \omega \cdot \Psi_{r d}\right]^{T}$
The matrices used are:

$$
\begin{aligned}
& A=\left[\begin{array}{ll}
a & b \\
c & d
\end{array}\right], B=\left[\begin{array}{ll}
0 & 0 \\
0 & 1
\end{array}\right] \\
& C=\left[\begin{array}{ll}
c_{1} & c_{2} \\
c_{3} & c_{4}
\end{array}\right], D=\left[\begin{array}{cc}
0 & -1 \\
1 & 0
\end{array}\right]
\end{aligned}
$$

Elements of matrix $A$ are:
$a=-\frac{R_{s} L_{r}}{L_{s} \cdot L_{r}-L_{m}^{2}}, b=\frac{R_{s} L_{m}}{L_{s} \cdot L_{r}-L_{m}^{2}}$
$c=\frac{R_{r} L_{m}}{L_{s} \cdot L_{r}-L_{m}^{2}}, d=-\frac{R_{r} L_{s}}{L_{s} \cdot L_{r}-L_{m}^{2}}$
Elements of matrix $C$ are:
$c_{1}=\frac{a}{R_{s} \cdot L_{r}}, c_{2}=-\frac{b}{R_{s} \cdot L_{r}}$
$c_{3}=\frac{c}{R_{s} \cdot L_{r}}, c_{4}=-\frac{d}{R_{s} \cdot L_{r}}$

Constant element's used are:

$$
k=\frac{3}{2} \cdot \frac{L_{m}}{L_{s} \cdot L_{r}-L_{m}^{2}}, \alpha=\frac{E_{2}}{U_{1}}, a_{f}=\frac{1}{T}
$$

We write an interesting, compact matrix model with two symmetric matrices:

$$
\begin{align*}
& {\left[\begin{array}{l}
\dot{\Psi}_{d} \\
\dot{\Psi}_{q}
\end{array}\right]=\left[\begin{array}{ll}
A & 0 \\
0 & A
\end{array}\right] \cdot\left[\begin{array}{l}
\Psi_{d} \\
\Psi_{q}
\end{array}\right]+\left[\begin{array}{l}
u_{d a}-u_{f q} \\
u_{q a}-u_{f d}
\end{array}\right]}  \tag{1}\\
& {\left[\begin{array}{l}
i_{d} \\
i_{q}
\end{array}\right]=\left[\begin{array}{ll}
C & 0 \\
0 & C
\end{array}\right] \cdot\left[\begin{array}{l}
\Psi_{d} \\
\Psi_{q}
\end{array}\right]}  \tag{2}\\
& M=k \cdot \Psi_{d}^{T} \cdot Q \cdot \Psi_{q}  \tag{3}\\
& \dot{\omega}=\frac{M-M_{s}}{J} \tag{4}
\end{align*}
$$

Equations (1), (2) look like state model and (3) and (4) are scalar depending of the previously terms.

In Fig. 1 is depicted the comprehensive model based on equations above.


Fig. 1 The comprehensive induction motor mathematical model

Fig. 2 presents the simulation results for speed, torque and feedback voltages.

Fluxes depend by sine and cosine functions. The speed can be approximated more accuratly using Berstain polynoms or by asymptotic relation.

## 2 Problem Formulation

We aproximate the feedback voltages as follows:

$$
\begin{equation*}
u_{f d}=-\alpha \cdot U_{s m} \cdot \frac{\omega_{s}}{\omega_{0}} \cdot\left(1-e^{-a_{f} \cdot t}\right) \cdot \cos \omega_{0} \cdot t \tag{5}
\end{equation*}
$$

$$
\begin{equation*}
u_{f q}=-\alpha \cdot U_{s m} \cdot \frac{\omega_{s}}{\omega_{0}} \cdot\left(1-e^{-a_{f} \cdot t}\right) \cdot \sin \omega_{0} \cdot t \tag{6}
\end{equation*}
$$

In the equations above $U_{s m}, \omega_{0}$ are the voltages supply parameters, $\omega_{s}$ is the steady state speed, $\alpha$ and $a_{f}$ are the calibration parameters

The nonlinear block N was replaced by Feedback Voltage Generator (FVG) and the mathematically model became almost linear, excepting torque equation.

The new simulation results depicted by Fig. 2 are close to those of nonlinear model..


Fig. 2 The nonlinear model and simulation results
Differences exist only on the start of the machine because the nonlinearities influences on start are relevant.


Fig. 3 The approximated linear model and simulation results

In this section we propose and solve a mathematical model based on transfer matrices for complex

## 3 Problem Solution

systems behavioral response. It's a complex subject to solve the linear matrix equations (1).

The general solution for the equation is depicted by the equations below.

The used eigenvalues $S_{1}, s_{2}$ are:

$$
\begin{aligned}
& s_{1}=\frac{(a+d)+\sqrt{(a-d)^{2}+4 b c}}{2} \\
& s_{2}=\frac{(a+d)-\sqrt{(a-d)^{2}+4 b c}}{2}
\end{aligned}
$$

The $s_{1}, s_{2}$ are the eigenvalues of matrix $A$, the $A_{0}, B_{0}, C_{0}, D_{0}, A_{i}(i=1,6), B_{i}(i=1,6), C_{i}(i=1,4)$ are
the constants depending of A matrix elements and feedback voltage parameters.

The $\phi_{i}(i=1,8), \alpha_{i}(i=1,8)$ depend of $A_{i}, B_{i}, C_{i}$.
The next step is to calculate the torque using the relation:

$$
M=k \cdot\left[\Psi_{r d} \cdot \Psi_{s q}-\Psi_{r q} \cdot \Psi_{s d}\right]
$$

The full expression for the torque is presented by the following equation

$$
\begin{array}{rlr}
M= & M_{0}+M_{1} \cdot e^{-a_{f} t}-M_{2} \cdot e^{-2 a_{f} t}+M_{2} \cdot e^{-2 a_{f} t}+M_{3} \cdot e^{-2 s_{1} t}+M_{4} \cdot e^{2 s_{2} t}+M_{5} \cdot e^{\left(s_{1}+s_{2}\right) t}+ \\
& +e^{s_{1} t} \cdot M_{6} \cdot \sin \left(\omega_{0} t-d_{3}-d_{5}\right)-e^{s_{1} t} \cdot M_{7} \cdot \sin \left(\omega_{0} t-d_{1}-d_{0}\right)-e^{s_{2} t} \cdot M_{8} \cdot \sin \left(\omega_{0} t-d_{3}+d_{7}\right)+ \\
& +e^{s_{2} t} \cdot M_{9} \cdot \sin \left(\omega_{0} t-d_{1}-d_{8}\right)+\left[-M_{10} \cdot \sin \left(\omega_{0} t-d_{4}-d_{5}\right)-M_{11} \cdot \cos \left(\omega_{0} t-\alpha_{2}-\alpha_{0}\right)\right] e^{\left(s_{1}-a_{f}\right) t}+ \\
& +\left[M_{12} \cdot \cos \left(\omega_{0} t-d_{4}+d_{7}\right)+M_{13} \cdot \cos \left(\omega_{0} t+d_{2}-d_{8}\right)\right] \cdot e^{\left(s_{2}-a_{f}\right) t} \\
& M_{0}=k \cdot \phi_{1} \cdot \phi_{3} \cdot \sin \left(\alpha_{3}-\alpha_{1}\right) & M_{6}=k \cdot \phi_{3} \cdot \phi_{5} \\
M_{1}= & k \cdot \phi_{2} \cdot \phi_{3} \cdot \cos \left(\alpha_{2}+\alpha_{3}\right)+ & M_{7}=k \cdot \phi_{1} \cdot \phi_{8} \\
& +k \cdot \phi_{1} \cdot \phi_{2} \cdot \sin \left(\alpha_{1}-\alpha_{4}\right) & M_{8}=k \cdot \phi_{3} \cdot \phi_{7} \\
M_{2}=k \cdot \phi_{2} \cdot \phi_{4} \cdot \cos \left(\alpha_{2}+\alpha_{4}\right) & M_{9}=k \cdot \phi_{1} \cdot \phi_{8} \\
M_{3}=k\left[B_{3} \cdot C_{3}-C_{3} \cdot A_{3}\right] & M_{10}=k \cdot \phi_{4} \cdot \phi_{5} \\
M_{4}=k\left[B_{4} \cdot C_{2}-C_{4} \cdot A_{4}\right] & M_{11}=k \cdot \phi_{2} \cdot \phi_{6} \\
M_{5}= & k \cdot B_{3} \cdot C_{2}+k \cdot B_{4} \cdot C_{1}- & M_{12}=k \cdot \phi_{4} \cdot \phi_{7} \\
& -k \cdot C_{3} \cdot A_{4}-k \cdot C_{4} \cdot A_{3} & M_{13}=k \cdot \phi_{2} \cdot \phi_{8}
\end{array}
$$

Analyzing the magnitude of $M_{i}(i=1,13)$ six terms can be neglected without losing the accuracy
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$$
\begin{equation*}
M=M_{0}-M_{1} \cdot e^{-a_{f} t}-M_{2} \cdot e^{-2 a_{f} t}-\left[M_{10} \cdot \sin \left(\omega_{0} t-\alpha_{4}-\alpha_{5}\right)+M_{11} \cdot \cos \left(\omega_{0} t+\alpha_{2}-\alpha_{6}\right)\right] \cdot e^{\left(s_{1}-a_{f}\right) t} \tag{8}
\end{equation*}
$$

$$
\begin{aligned}
& {\left[\begin{array}{l}
\Psi_{d} \\
\Psi_{q}
\end{array}\right]=\left[\begin{array}{cc}
e^{A t} & 0 \\
0 & e^{A t}
\end{array}\right] \cdot\left[\begin{array}{l}
\Psi_{d}(0) \\
\Psi_{q}(0)
\end{array}\right]+\int_{0}^{t} e^{A(t-\tau)} \cdot\left[\begin{array}{l}
u_{d a}(\tau)-u_{f q}(\tau) \\
u_{q a}(\tau)+u_{f d}(\tau)
\end{array}\right] \cdot d \tau} \\
& e^{A t}=L^{-1}\left[(s I-A)^{-1}\right]=\frac{1}{s_{2}-s_{1}}\left[\begin{array}{cc}
\left(d-s_{1}\right) e^{s_{1} t}-\left(d-s_{2}\right) e^{s_{2} t} & -b e^{s_{1} t}+b e^{s_{2} t} \\
-c e^{s_{1} t}+c e^{s_{2} t} & \left(a-s_{1}\right) e^{s_{1} t}-\left(a-s_{2}\right) e^{s_{2} t}
\end{array}\right] \\
& {\left[\begin{array}{l}
\Psi_{d} \\
\Psi_{q}
\end{array}\right]=\left[\begin{array}{l}
\Psi_{s d} \\
\Psi_{r d} \\
\Psi_{s q} \\
\Psi_{r q}
\end{array}\right]=\frac{1}{s_{2}-s_{1}}\left\{\left[\begin{array}{cccc}
A_{0} e^{s_{t} t}-B_{0}{ }^{s_{2} t} & -b e^{s_{1} t}+b e^{s_{2} t} & 0 & 0 \\
-c e^{s_{1} t}+c e^{s_{2} t} & C_{0} e^{s_{1}}-D_{0} e^{s_{2} t} & 0 & 0 \\
0 & 0 & A_{0} e^{s_{1} t}-B_{0} e^{s_{2} t} & -b e^{s_{1} t}+b e^{s_{2} t} \\
0 & 0 & -c e^{s_{1 t}}+c e^{s_{2} t} & C_{0} e^{s_{1} t}-D_{0} e^{s_{2} t}
\end{array}\right] \cdot\left[\begin{array}{l}
\Psi_{s d}(0) \\
\Psi_{r d}(0) \\
\Psi_{s q}(0) \\
\Psi_{r q}(0)
\end{array}\right]+\right.} \\
& \left.+\left[\begin{array}{l}
-\phi_{1} \cos \left(\omega_{0} t-\alpha_{1}\right)+\phi_{2} \sin \left(\omega_{0} t-\alpha_{2}\right) e^{-a_{f} t}+A_{3} e^{s_{1} t}+A_{4} e^{s_{2} t} \\
-\phi_{3} \cos \left(\omega_{0} t-\alpha_{3}\right)+\phi_{4} \cos \left(\omega_{0} t-\alpha_{4}\right) e^{-a_{f} t}+B_{3} e^{s_{1} t}+B_{4} e^{s_{2} t} \\
-\phi_{1} \sin \left(\omega_{0} t-\alpha_{1}\right)-\phi_{2} \cos \left(\omega_{0} t+\alpha_{2}\right) e^{-a_{f} t}+C_{1} e^{s_{1} t}+C_{2} e^{s_{2} t} \\
-\phi_{3} \sin \left(\omega_{0} t-\alpha_{3}\right)+\phi_{4} \cos \left(\omega_{0} t-\alpha_{4}\right) e^{-a_{f} t}+C_{3} e^{s_{1} t}+C_{4} e^{s_{2} t}
\end{array}\right]\right\}
\end{aligned}
$$

The speed is determined from the equation:

$$
\begin{align*}
\frac{d \omega}{d t}= & \frac{\left(M-M_{s}\right)}{J} \text {, where M is expressed by equation (8). } \\
\omega= & \frac{M_{0}-M_{s}}{J} t+\frac{M_{1}}{J a_{f}}\left(1-e^{-a_{f} t}\right)-\frac{M_{2}}{2 J a_{f}}\left(1-e^{-2 a_{f} t}\right)-\frac{M_{10}\left(s_{1}-a_{f}\right)}{J\left[\left(s_{1}-a_{f}\right)^{2}+\omega_{0}^{2}\right]} e^{\left(s_{1}-a_{f}\right) t} \sin \left(\omega_{0} t-\alpha_{4}-\alpha_{5}\right)- \\
& -\frac{M_{10}\left(s_{1}-a_{f}\right)}{J\left[\left(s_{1}-a_{f}\right)^{2}+\omega_{0}^{2}\right]^{2}} \sin \left(\alpha_{4}+\alpha_{5}\right)+\frac{M_{10} \omega_{0}}{J\left[\left(s_{1}-a_{f}\right)^{2}+\omega_{0}^{2}\right]} e^{\left(s_{1}-a_{f}\right) t} \cos \left(\omega_{0} t-\alpha_{4}-\alpha_{5}\right)- \\
& -\frac{M_{10} \omega_{0}}{J\left[\left(s_{1}-a_{f}\right)^{2}+\omega_{0}^{2}\right]^{2}} \cos \left(\alpha_{4}+\alpha_{5}\right)-\frac{M_{11}\left(s_{1}-a_{f}\right) \omega_{0}}{J\left[\left(s_{1}-a_{f}\right)^{2}+\omega_{0}^{2}\right]^{\left(s_{1}-a_{f}\right) t}} \cos \left(\omega_{0} t+\alpha_{2}-\alpha_{6}\right)+  \tag{9}\\
& \left.+\frac{M_{11}\left(s_{1}-a_{f}\right) \omega_{0}}{J\left[\left(s_{1}-a_{f}\right)^{2}+\omega_{0}^{2}\right.}\right]^{\cos \left(\alpha_{2}+\alpha 6\right)-\frac{M_{11} \cdot \omega_{0}}{J\left[\left(s_{1}-a_{f}\right)^{2}+\omega_{0}^{2}\right.} e^{\left(s_{1}-a_{f}\right) t} \sin \left(\omega_{0} t+\alpha_{2}-\alpha_{6}\right)+} \\
& +\frac{M_{11} \cdot \omega_{0}}{J\left[\left(s_{1}-a_{f}\right)^{2}+\omega_{0}^{2}\right]^{2}} \sin \left(\alpha_{2}-\alpha 6\right)
\end{align*}
$$

The $M_{0} \approx M_{s}$ for stability of the drive and the last three terms are negligible because of the high value of $\omega_{0}^{2}$ from the denominators. The simplified form of the speed is:

$$
\omega_{0}=\frac{M_{1}}{J \cdot a_{f}}\left(1-e^{-a_{f} t}\right)-\frac{M_{2}}{2 J a_{f}}\left(1-e^{-2 a_{f} t}\right)
$$

The most important thing is the possibility to determine the transfer matrix from voltage to fluxes.

$$
\begin{aligned}
& \Psi_{d}=\left[\begin{array}{l}
\Psi_{s d} \\
\Psi_{r d}
\end{array}\right]=[s I-A]^{-1} \cdot\left[\begin{array}{l}
u_{s d} \\
u_{r d}
\end{array}\right] \\
& \Psi_{d}=[s I-A]^{-1} \cdot\left[\begin{array}{c}
u_{s m} \cdot \sin \left(\omega_{0} t\right) \\
-\left(1-e^{-a_{f} t}\right) \cdot u_{r m} \cdot \sin \left(\omega_{0} t\right)
\end{array}\right] \\
& L\left\{\left(1-e^{-a_{f} t}\right) u_{r m} \sin \left(\omega_{0} t\right)\right\}=\frac{-2 \cos a_{f} \cdot s}{k_{0} \omega_{0}\left[\left(s+a_{f}\right)^{2}+\omega_{0}^{2}\right]} \\
& u_{r m}=\frac{\omega_{s}}{\omega_{0}} \cdot u_{s m}
\end{aligned}
$$

Based on the equations above we get the following transfer matrix:

$$
\begin{aligned}
& {\left[\begin{array}{l}
\Psi_{d} \\
\Psi_{q}
\end{array}\right]=\left[\begin{array}{l}
G_{d}(s) \\
G_{q}(s)
\end{array}\right] \cdot u_{s m}, \text { where }} \\
& G_{d}(s)=\left[\begin{array}{ll}
G_{1} & G_{2} \cdot G_{5} \\
G_{3} & G_{4} \cdot G_{5}
\end{array}\right] \\
& G_{q}(s)=\left[\begin{array}{ll}
G_{1} & G_{2} \cdot G_{6} \\
G_{3} & G_{4} \cdot G_{6}
\end{array}\right]
\end{aligned}
$$

Each of the $G_{i}(i=1 . .6)$ terms has an expression depending of the electric and the mechanic parameters.

The parameters from a real induction motor were used for calculating the six terms of the transfer matrix. The values of the motor parameters are depicted in Fig. 3 and the resulting values of transfer matrix terms are displayed by the equation set (10).

$$
\begin{align*}
& G_{1}(s)=\frac{s+667}{s^{2}+1298 s+8250} \\
& G_{2}(s)=\frac{875}{s^{2}+1298 s+8250} \\
& G_{3}(s)=\frac{467}{s^{2}+1298 s+8250}  \tag{10}\\
& G_{4}(s)=\frac{s+625}{s^{2}+1298 s+8250} \\
& G_{5}(s)=\frac{5.58 s+7.784}{s^{2}+5.58 s+98604} \\
& G_{6}(s)=\frac{2.79 s^{2}+7.784 s-275083}{s^{3}+5.58 s^{2}+98604 s}
\end{align*}
$$

Based on the parameters and transfer matrix a Matlab Simulink simulation model was created in order to validate the proposed mathematical model.

As can be observed in Fig. 4 the curves obtained for torque and speed match very accurate the ones obtained from the nonlinear model and linear approximated model [10].

This result allows the implementation of the mathematical model into an intelligent electric inverter and to drive an AC induction motor through all functioning regimes to confirm the results.


Fig. 4 Induction motor: a) transfer matrices model; b) simulation results.

## 4 Conclusion

An original approach based on the asymptotic linearization feedback voltages and the system theory was developed in this paper in order to solve the mathematical model of the induction machine.

To ease finding the solution for the proposed model, an interesting matrix system was constructed in order to accurately approximate the system response.

Then the solution equations were modeled and simulated, and the results were compared with the ones of the consecrated model with synthesized feedback voltages.

After determining the solutions for field components the expressions for torque, speed and transfer matrix on channel voltage-field were calculated and then implemented into a simulation model for results analysis.

The results are original and very close with the ones from the nonlinear model validating the method developed here.

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