# A numerical study of natural convection in a cavity with two baffles attached to its vertical walls 

M. GHASSEMI<br>Department of Mechanical Engineering<br>K.N.Toosi University of<br>Technology<br>I. R. IRAN

M. PIRMOHAMMADI<br>Department of Mechanical Engineering<br>K.N.Toosi University of<br>Technology<br>I. R. IRAN

GH. A. SHEIKHZADEH<br>Department of Mechanical<br>Engineering<br>University of Kashan<br>I. R. IRAN


#### Abstract

The purpose of this study is to investigate the effect of two insulated horizontal baffles placed at the walls of a differentially heated square cavity. The vertical walls are at different temperatures while the horizontal walls are adiabatic. In our formulation of governing equations, mass, momentum and the energy equations are applied to the cavity and the baffles. To solve the governing differential equations a finite volume code based on Pantenkar's simpler method is utilized. The Results are presented for Rayleigh number from $10^{4}$ up to $10^{7}$ and are in form of streamlines, isotherms as well as Nusselt number. It is observed that the two baffles trap some fluid in the cavity and affect the flow fields. Also, the Nusselt number increases as Rayleigh number increases and decreases with baffle length. Finally it is shown that Nusselt number changes with baffles position.


Key-words: -Natural Convection, Cavity, baffles, Nusselt Number

## 1 Introduction

The phenomenon of natural convection heat transfer plays an important role, both in nature and in engineering systems. Many investigations have been performed for cavities both theoretically and experimentally for a wide range of Ra (Rayleigh number) [1-2]. Natural convection in an air filled cavity with vertical walls that are heated and cooled while its horizontal walls are adiabatic has received a great consideration because many of the industrial applications employ this concept as a prototype. Noticeable examples include heating and ventilation of rooms, solar collector systems, and electronic cooling devices. In many applications, for some reasons, attaching baffle(s) to its vertical wall or to its horizontal wall(s) partitions the cavity. Recently studies of heat transfer and fluid flow characteristics of partitioned cavity have come under scrutiny both numerically and experimentally [3-11].
Bajorek and Llyod [3] studied experimentally a differentially heated air filled square cavity for various Ra from $1.25 \times 10^{5}$ to $2.16 \times 10^{6}$. The insulated baffle is attached to the horizontal walls at positions in the middle. Non-dimensional baffle length is 0.25 . The effect of the baffle positions was not considered. It was found that the baffles significantly influence the heat transfer rate and the average Nusselt numbers reduced to approximately $15 \%$ compared to
the non-partitioned cavity. Jetli et al [6] studied numerically a differentially heated air filled square partitioned cavity for various Ra from $10^{4}$ to $3.55 \times 10^{5}$ with three different combinations of the baffle positions. Non-dimensional baffles length is fixed at 0.33 . The results clearly demonstrate that the baffles positions have a significant effect on the heat transfer and flow characteristics of the fluid. For all baffles locations, the average Nusselt number is smaller than the corresponding value in a cavity without baffle. Ambarita studied numerically a differentially heated square cavity with two perfectly insulated baffles were attached to its horizontal walls at symmetric position [9]. A parametric study was carried out using following parameters: Rayleigh number from $10^{4}$ to $10^{8}$, non-dimensional thin baffles length were $0.6,0.7,0.8$, non-dimensional baffle position were 0.2 to 0.8 .
The purpose of this study is to investigate the effect of two insulated horizontal baffles placed at the walls of a differentially heated square cavity. The vertical walls are at different temperatures while the horizontal walls are adiabatic.

## 2 Problem definition

The typical cavity with the boundaries and its coordinates system are depicted in Fig.1. In this
study only the square cavity is considered, H denotes its height and width. The cavity is differentially heated, left and right walls are isothermal at $\mathrm{T}_{\mathrm{h}}$ and $\mathrm{T}_{\mathrm{c}}$ respectively ( $\mathrm{T}_{\mathrm{h}}>\mathrm{T}_{\mathrm{c}}$ ) and horizontal walls are adiabatic.
Two thin baffles with non-dimensional length $\mathrm{L}_{\mathrm{b}}$, perfectly insulated, are attached to the left and right wall. The non-dimensional position of the left baffle from the bottom wall and right baffle from the top wall are the same and denoted by $\mathrm{D}_{\mathrm{b}}$.


Fig. 1 Schematic of the square cavity with thin insulated baffles attached to the vertical walls

The governing equations are converted into the nondimensional form by defining the non-dimensional variables:
$X=\frac{x}{H}, Y=\frac{y}{H}, U=\frac{u H}{\alpha}, V=\frac{v H}{\alpha}, P=\frac{p H^{2}}{\rho \alpha^{2}}, \theta=\frac{T-T_{c}}{T_{h}-T_{c}}$
Based on these non-dimensional variables, the governing equations are obtained as follows:
$\frac{\partial U}{\partial X}+\frac{\partial V}{\partial Y}=0$
$U \frac{\partial U}{\partial X}+V \frac{\partial U}{\partial Y}=-\frac{\partial P}{\partial X}+\operatorname{Pr}\left(\frac{\partial^{2} U}{\partial X^{2}}+\frac{\partial^{2} U}{\partial Y^{2}}\right)$
$U \frac{\partial V}{\partial X}+V \frac{\partial V}{\partial Y}=-\frac{\partial P}{\partial X}+\operatorname{Pr}\left(\frac{\partial^{2} V}{\partial X^{2}}+\frac{\partial^{2} V}{\partial Y^{2}}\right)+\operatorname{Ra} \operatorname{Pr} \theta$
$U \frac{\partial \theta}{\partial X}+V \frac{\partial \theta}{\partial Y}=\left(\frac{\partial^{2} \theta}{\partial X^{2}}+\frac{\partial^{2} \theta}{\partial Y^{2}}\right)$
The boundary conditions are:
On the left wall: $U=V=0, \theta=1$
On the right wall: $U=V=0, \theta=0$
On the top and bottom walls: $U=V=0, \partial \theta / \partial Y=0(8)$
On the baffles: $U=V=0, \partial \theta / \partial Y=0$
In order to compare total heat transfer rate, Nusselt number is used. The local and average Nusselt numbers are defined as follows:

$$
\begin{equation*}
N u_{y}=-\frac{\left.\partial \theta \theta_{\partial X}\right|_{X=0}}{\left(\theta_{h}-\theta_{c}\right)}, \quad \bar{N} u=\int_{0}^{1} N u_{y} d Y \tag{10}
\end{equation*}
$$

## 3 Numerical procedure

The coupled governing equations are transformed into sets of algebraic equations using finite volume
method and are solved by the SIMPLER algorithm [15]. The staggered grid system is used and the convective terms are handled by the power law scheme.

### 3.1 Validation of the numerical code

In order to make sure that the developed codes are free of error coding, a validation test is conducted. Calculations for an air filled cavity without baffle for $\mathrm{Ra}=10^{4}$ to $10^{7}$ are carried out and the results are shown in Table 1. The results of the previous publications for the same problem are also presented in the Table 1. Data from the table shows that the results of the code, even though there are some differences, do agree very well with the previous results. Based on this successful validation, the problem is solved by using the code.

Table 1 Comparison of the present result and the

| Reference | Average Nusselt numbers, $\bar{N} u$ |  |  |  |
| :--- | :---: | :---: | :---: | :---: |
|  | $\mathrm{Ra}=10^{4}$ | $\mathrm{Ra}=10^{5}$ | $\mathrm{Ra}=10^{6}$ | $\mathrm{Ra}=10^{7}$ |
|  | 2.4468 | 5.5231 | 8.8359 | - |
|  | 2.244 | 4.5236 | 8.8554 | - |
| Nag[11] | 2.24 | 4.51 | 8.82 | - |
| Shi[1] | 2.247 | 4.532 | 8.893 | 16.935 |
| Bilgen[8] | 2.245 | 4.521 | 8.8 | 16.629 |
| Ambarita[9] | 2.228 | 4.514 | 8.804 | 16.52 |
| Present | 2.25 | 4.53 | 8.85 | 16.43 |

Flow and temperature fields for the square cavity without baffle are presented in Fig.2. It is observed that as natural convection strengthened, temperature contours show slight deviation from the pure conduction case with the isotherms becoming skewed. Under high Ra conditions, the degree of distortion from the pure conduction case is varying marked and the contour lines become almost horizontal lines around the center of the cavity; Fig.2(a). The rise of the fluid due to heating on the left wall and consequent falling of the fluid on the right wall creates a clockwise rotating vortex, referred to as the primary vortex; Fig.2(b).
Another feature of these streamlines patterns is that the streamlines become more packed next to the sidewall as the Ra increases. This suggests that the flow move faster as natural convection is intensified. The maximum absolute value of stream function can be viewed as a measure of the intensity of natural convection in the cavity. As the Ra increases, the maximum absolute value of the stream function increases.
(a)

(b)

##  <br>  <br> $|\psi|_{\max }=5.035$

,

$|\psi|_{\max }=9.7$





$|\psi|_{\max }=30.33$
Fig. 2 Isotherms (a) and Streamlines (b) of natural convection in a square cavity

## 4 Results and discussion

In order to understand the flow and temperature fields and heat transfer characteristics of the typical cavity a total of 64 cases are considered. To study the effects of baffle position and baffle length, the non-dimensional baffle positions $\mathrm{D}_{\mathrm{b}}=0.2,0.4,0.6$ and 0.8 and the non-dimensional baffle lengths $\mathrm{L}_{\mathrm{b}}=0.5,0.6,0.7$ and 0.8 are considered. The fluid inside the cavity is dry air with $\operatorname{Pr}=0.7$ and Rayleigh number varied from $10^{4}$ to $10^{7}$.

### 4.1 Flow and temperature fields

Flow fields for the typical cases with nondimensional baffle length $\mathrm{L}_{\mathrm{b}}=0.6$ are presented in Fig.3. The plots are arranged going from left to right
with the ascending of Ra and from top to bottom with the ascending of $D_{b}$.
A glance at Fig. 3 shows that the flow fields can be divided into two different patterns. The first pattern is the fluid circulates and creates a large primary vortex strangled by the baffles. The second pattern is fluid separated into two different vortexes. The two baffles create fluid trapping phenomena in the cavity. The first raw of Fig. 3 is the flow fields for the case when $\mathrm{D}_{\mathrm{b}}=0.2$. At a low Rayleigh number $\mathrm{Ra}=10^{4}$ there are two trapped fluids in the cavity, bottom trapped fluid and top trapped fluid. The bottomtrapped fluid exists between bottom baffle and bottom adiabatic wall and top trapped fluid exists between top baffle and top adiabatic wall. This is because at $\mathrm{Ra}=10^{4}$ natural convection is too weak to make the trapped fluids moving and also the space of trapped fluids is limited due to a small $D_{b \text {. The }}$ natural convection creates a primary vortex strangled by these two trapped fluids. By increasing Ra the primary vortex is strong enough to penetrate the two trapped fluid.
The second raw of Fig. 3 is the flow field when $\mathrm{D}_{\mathrm{b}}=0.4$. The figure shows that even for a low Rayleigh number $\mathrm{Ra}=10^{4}$ the primary vortex divide into two vortexes. This is because the space between each baffle and the nearest horizontal wall is larger compared to the corresponding cases when $\mathrm{D}_{\mathrm{b}}=0.2$.
In order to satisfy continuity, the circulation on the bottom side is separated from topside. At low Ra a weak vortex is formed between two baffles. By increasing Ra these vortexes will be strengthened.
The flow field for case with $D_{b}=0.6$ are shown in the third column of the fig3. At $\mathrm{Ra}=10^{4}$ there are two vortexes in the cavity and a trapped fluid exists between the two baffles. The fourth raw of Fig. 3 is the flow fields for the case when $D_{b}=0.8$. In this case by increasing space between two baffles and reducing the space between baffles and adiabatic walls, two separated vortexes will be merged and again primary vortex between baffles and twotrapped fluid will be formed. At a low $\mathrm{Ra}\left(\mathrm{Ra}=10^{4}\right)$ there are two trapped fluids in the cavity, bottom trapped fluid and top trapped fluid. The bottomtrapped fluid exists between bottom baffle and bottom adiabatic wall and top trapped fluid exists between top baffle and top adiabatic wall. This is because at $\mathrm{Ra}=10^{4}$ natural convection is too weak to make the trapped fluids moving. By increasing Ra the primary vortex tends to penetrate the trapped fluid and trapped vortex will be strengthened.
Temperature fields when $\mathrm{L}_{\mathrm{b}}=0.6$ are presented in Fig. 4.


Fig. 3 Streamlines for $L_{b}=0.6$ and various $R a$ and $D_{b}$

The plots are arranged going from left to right with the ascending of Ra and from top to bottom with the ascending of $\mathrm{D}_{\mathrm{b}}$. Contour level increments for each case are kept constant at 0.1 . Indeed in comparison to the square cavity without baffles the appearance of the temperature fields is strongly modified due to the presence of the two insulated baffles.
The first raw of Fig. 4 is for $D_{b}=0.2$. At $\mathrm{Ra}=10^{4}$ when the natural convection is weak and in the region where fluid is trapped the conduction heat transfer only occurred. In the region where natural convection takes place the isotherms show deviations from the pure conduction case with the contour lines becoming skewed. In the trapped fluid areas heat transfer is inactive due to presence of the insulated baffles and stagnant fluids. At $\mathrm{Ra}=10^{5}$ the isotherms
in the top left and bottom right of the cavity become more skewed but almost vertical between the two baffles and insulated walls. This is because natural convection is more vigorous in the top left and bottom right of the cavity but becomes weak between the two baffles and insulated walls. For $\mathrm{Ra}=10^{6}$, primary vortex is strong enough to penetrate the two trapped fluid then Isotherms in these regions are skewed and convection heat transfer is appearance.
The second raw of Fig. 4 is the temperature fields for cases when baffle positions $\mathrm{D}_{\mathrm{b}}=0.4$. The figure shows that in the regions between baffles and adiabatic walls the isotherms are skewed and conduction heat transfer is less compared to the corresponding cases when $\mathrm{D}_{\mathrm{b}}=0.2$.


Fig. 4 Isotherms for $L_{b}=0.6$ and various $R a$ and $D_{b}$

By increasing the Ra the isotherms are more skewed in these regions and convection heat transfer is clearer. The third raw is for case when $\mathrm{D}_{\mathrm{b}}=0.6$. At $\mathrm{Ra}=10^{4}$ where the flow is still separated into two vortexes and a trapped fluid exists between the two baffles the isotherms are skewed on the vortex areas but almost vertical in the trapped fluid area. At $\mathrm{Ra}=10^{5}$ the isotherms are more skewed in the vortex areas but less vertical in the trapped fluid area since intensity of the natural Convection increases. At $\mathrm{Ra}=10^{6}$ two vortexes have strengthened.
Further there is no trapped fluid between baffles. The forth raw of Fig. 4 are the temperature fields for cases when the baffle positions $D_{b}=0.8$. The figure shows that the Isotherms lines in the regions between baffles and adiabatic walls are vertical and only conduction heat transfer occurred there. And by increasing the Rayleigh number isotherms are skewed and convection heat transfer is clearer.

### 4.2 Heat transfer

In order to evaluate how the presence of the two baffles affects the heat transfer rate through the cavity, the average Nusselt number will be discussed. The average Nusselt numbers for all cases as a function of Rayleigh number and for various parameters, $\mathrm{L}_{\mathrm{b}}=0.5, \mathrm{~L}_{\mathrm{b}}=0.6$, and $\mathrm{L}_{\mathrm{b}}=0.7, \mathrm{~L}_{\mathrm{b}}=0.8$ are presented in Table 2.
Generally, as expected $\bar{N} u$ is an increasing function of Ra and a decreasing function of $\mathrm{L}_{\mathrm{b}}$. Also the Nusselt number is baffle position dependent. As the Ra increases the maximum absolute value of the stream function increases. This means the intensity of natural convection in the cavity increases as well as the temperature gradient near the isotherm walls. Table 2 shows the case with $\mathrm{L}_{\mathrm{b}}=0.6$ and For $\mathrm{D}_{\mathrm{b}}=$ $0.2, \bar{N} u$ is decreased by $32 \%$ at $\mathrm{Ra}=10^{4}$ and is $12 \%$ at $\mathrm{Ra}=10^{7}$, the average is $22 \%$. For $\mathrm{D}_{\mathrm{b}}=0.4, \mathrm{Nu}$ is decreased by $50 \%$ at $\mathrm{Ra}=10^{4}$ and is $9 \%$ at $\mathrm{Ra}=10^{7}$, the average is $29.5 \%$. These values suggest that the
case with $D_{b}=0.4$ blocks the heat transfer rate more effectively than the case with $D_{b}=0.2$. For $D_{b}=0.6$, $\bar{N} u$ is decreased by $47 \%$ at Ra $=10^{4}, 13 \%$ at $\mathrm{Ra}=10^{7}$, the average is $30 \%$. For $\mathrm{D}_{\mathrm{b}}=0.8, \bar{N} u$ is decreased by $35 \%$ at $\mathrm{Ra}=10^{4}$ and is $14 \%$ at $\mathrm{Ra}=10^{7}$, the average is $24.5 \%$. The results clearly demonstrate that the baffles positions have a significant effect on the heat transfer and flow characteristics of the fluid. For all baffles locations, the average Nusselt number is smaller than the corresponding value in a cavity without baffle.

Table 2 Nusselt number for various Ra, $\mathrm{L}_{\mathrm{b}}$ and $\mathrm{D}_{\mathrm{b}}$

|  | $\mathrm{Ra}=10^{4}$ | $\mathrm{Ra}=10^{5}$ | $\mathrm{Ra}=10^{6}$ | $\mathrm{Ra}=10^{7}$ |
| :---: | :---: | :---: | :---: | :---: |
| $\mathrm{D}_{\mathrm{b}}=0.2$ | 1.615 | 3.39 | 7.12 | 14.66 |
| $\mathrm{D}_{\mathrm{b}}=0.4$ | 1.18 | 3.35 | 8.12 | 15.31 |
| $\mathrm{D}_{\mathrm{b}}=0.6$ | 1.22 | 3.21 | 7.8 | 14.29 |
| $\mathrm{D}_{\mathrm{b}}=0.8$ | 1.66 | 3.81 | 7.62 | 14.25 |
| $\mathbf{L}_{\mathrm{b}}=\mathbf{0 . 5}$ |  |  |  |  |
|  | $\mathrm{Ra}=10^{4}$ | $\mathrm{Ra}=10^{5}$ | $\mathrm{Ra}=10^{6}$ | $\mathrm{Ra}=10^{7}$ |
| $\mathrm{D}_{\mathrm{b}}=0.2$ | 1.53 | 3.33 | 6.96 | 14.4 |
| $\mathrm{D}_{\mathrm{b}}=0.4$ | 1.11 | 3.07 | 8.04 | 15.24 |
| $\mathrm{D}_{\mathrm{b}}=0.6$ | 1.19 | 3.19 | 7.63 | 14.15 |
| $\mathrm{D}_{\mathrm{b}}=0.8$ | 1.46 | 3.66 | 7.37 | 14.11 |
| $\mathrm{L}_{\mathrm{b}}=0.6$ |  |  |  |  |
|  | $\mathrm{Ra}=10^{4}$ | $\mathrm{Ra}=10^{5}$ | $\mathrm{Ra}=10^{6}$ | $\mathrm{Ra}=10^{7}$ |
| $\mathrm{D}_{\mathrm{b}}=0.2$ | 1.49 | 3.21 | 6.8 | 14.25 |
| $\mathrm{D}_{\mathrm{b}}=0.4$ | 1.07 | 2.71 | 7.9 | 15.15 |
| $\mathrm{D}_{\mathrm{b}}=0.6$ | 1.15 | 2.98 | 7.4 | 13.78 |
| $\mathrm{D}_{\mathrm{b}}=0.8$ | 1.34 | 3.54 | 7.26 | 14.07 |
| $\mathbf{L}_{\mathrm{b}}=\mathbf{0 . 7}$ |  |  |  |  |
|  | $\mathrm{Ra}=10^{4}$ | $\mathrm{Ra}=10^{5}$ | $\mathrm{Ra}=10^{6}$ | $\mathrm{Ra}=10^{7}$ |
| $\mathrm{D}_{\mathrm{b}}=0.2$ | 1.40 | 3.11 | 6.64 | 13.94 |
| $\mathrm{D}_{\mathrm{b}}=0.4$ | 1.05 | 2.55 | 7.7 | 15.12 |
| $\mathrm{D}_{\mathrm{b}}=0.6$ | 1.12 | 2.89 | 7.3 | 13.5 |
| $\mathrm{D}_{\mathrm{b}}=0.8$ | 1.26 | 3.06 | 7.14 | 14.02 |
| $\mathrm{L}_{\mathrm{b}}=0.8$ |  |  |  |  |

## 5 Conclusion

Heat transfer by natural convection in a differentially heated square cavity with two thin insulated baffles has been numerically studied. The cavity is performed by vertical isothermal walls and adiabatic horizontal walls. Two thin insulated baffles are attached to its vertical walls at symmetric positions. In order to understand the flow and temperature fields and heat transfer characteristics of the typical cavity a total of 64 cases are considered. To study the effects of the baffle position, the nondimensional baffle positions $D_{b}=0.2,0.4,0.6$ and 0.8 are considered. To study the effects of baffle length, the non-dimensional baffle lengths $\mathrm{L}_{\mathrm{b}}=0.5,0.6,0.7$ and 0.8 are considered and Ra varies from $10^{4}$ to $10^{7}$. Two different flow field patterns are observed. The first pattern is flow fields with two different vortexes separated by a trapped fluid between the baffles and the second pattern is flow field with a primary vortex strangled by two trapped fluids. The flow field
pattern is non-dimensional baffle positions and Rayleigh numbers dependent. It is also observed that the Nusselt number increases as Rayleigh number increases and decreases with baffle length

## References:

[1] X. Shi and J. M. Khodadadi, Laminar Natural Convection Heat Transfer in a Differentially Heated Square Cavity Due to a Thin Fin on the Hot Wall, Journal of Heat Transfer, 125, 2003, pp. 624-634.
[2] V. Mariani and I. Moura Belo, Numerical Studies of Natural Convection in a Square Cavity, Thermal Engineering, 5, 2006, pp. 68-72.
[3] S.M. Bajorek and J.R. Lloyd, Experimental Investigation of Natural Convection in Partitioned Enclosures, Journal of Heat Transfer, 104, 1982, pp.527-532.
[4] S.H. Tasnim and M. R. Collins, Numerical Analysis of Heat Transfer in a Square Cavity with a Baffle on the Hot Wall, Int. Comm. Heat Mass Transfer, 31, 2004, pp.639-650.
[5] F. Ampofo, Turbulent natural convection in an air filled partitioned square cavity, Int. J. of Heat and Fluid Flow, 25, 2004, pp. 103-114.
[6] R. Jetli, S. Acharya, and E. Zimmerman, Influence of Baffle Location on Natural Convection in a Partially Divided Enclosure, Numerical Heat Transfer,10,1986, pp.521-536.
[7] E.Bilgen, Natural Convection in Enclosures with Partial Partitions, Renewable Energy, 26, 2002, pp.257-270.
[8] E. Bilgen, Natural Convection in Cavities with a Thin Fin on the Hot Wall, Int. J. Heat and Mass Transfer, 48, 2005, pp.3493-3505.
[9] H. Ambarita, K. Kishinami, M.Dimaruya, T.Saitoh, H.Takahashi, and J.Suzuki, Laminar Natural Convection Heat Transfer in an Air Filled Square Cavity With Two Insulated Baffles Attached to its Horizontal Walls, Thermal Science \& Engineering, vol.14, 3, 2006, pp.35-46.
[10] M. Hortmann, M. Peric, and G Scheuerer, Volume Multigrid Prediction of Laminar Natural Convection; Bench-Mark Solutions, Int. J. Numerical Methods in Fluids, 11, 1990, pp. 187-207.
[11] A. Nag, A. Sarkar, and V. M.K. Sastri, Natural Convection in a Differentially Heated Square Cavity with a Horizontal Partition Plate on the Hot Wall, Comput. Methods Appl. Mech. Eng., 110, 1993, pp.143-156.
[12] S.V. Patankar, Numerical Heat Transfer and Fluid Flow, Hemisphere, Washington, DC, 1980.

